

GEM M-I Coupling Campaign Tutorials



# Electrodynamics of M-I Coupling

by

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# Outline



- Introduction: M-I Coupling Campaign WG2 (Electrodynamics of M-I Coupling)
- Magnetosphere – main electrodynamic driver of ionosphere, via **field-aligned currents**
- Magnetosphere-ionosphere link
  - **Parallel electric fields**
  - Particle energization; parallel electron beams; **aurora**
  - **Ionospheric conductivity**; Joule heating
- Time-dependent picture
  - **Alfvén waves**; ionospheric Alfvén resonator
  - Non-MHD effects; Wave-particle interactions
- Observations; FAST mission
- Issues in M-I Coupling Campaign WG2

# GEM M-I Coupling Campaign – WG2

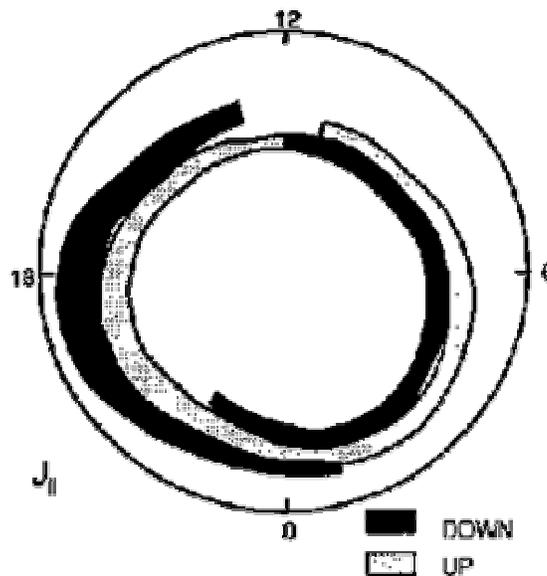


- WG1 – Ionospheric Plasma in the Magnetosphere
- **WG2 – Electrodynamics of M-I Coupling**
  - understanding how EM energy is transferred between the ionosphere and magnetosphere, and how several factors (e.g. field-aligned currents, electric fields, waves) affect this transfer at different scales
  - Connections with other campaigns (IM-Storms, GGCM) – **M-I coupling is just a part in the overall solar-wind-magnetosphere-ionosphere system**

# Driver: Field-aligned currents



- Linkage between solar wind-magnetosphere system and ionosphere – act as a **dynamo**
- **Region 1** and **region 2** currents; Mantle/NBZ currents at the ionosphere
- Magnetospheric formation: due to pressure gradients  $\nabla_{\perp} \cdot \mathbf{j}_{\perp} + \nabla_{\parallel} \cdot \mathbf{j}_{\parallel} = 0$
- FACs **close through the ionosphere**, in the auroral zones



- Ionosphere is **collisional**, so EM energy is dissipated there;  **$E_{\perp}$  electric fields must exist**, and so **the ionosphere is a load for the current system**

# Ionospheric conductivity: Hall and Pedersen currents



- Ionosphere – one-fluid model not applicable, need to use three (or more)-fluid model
- Relation between current and **E**-field:

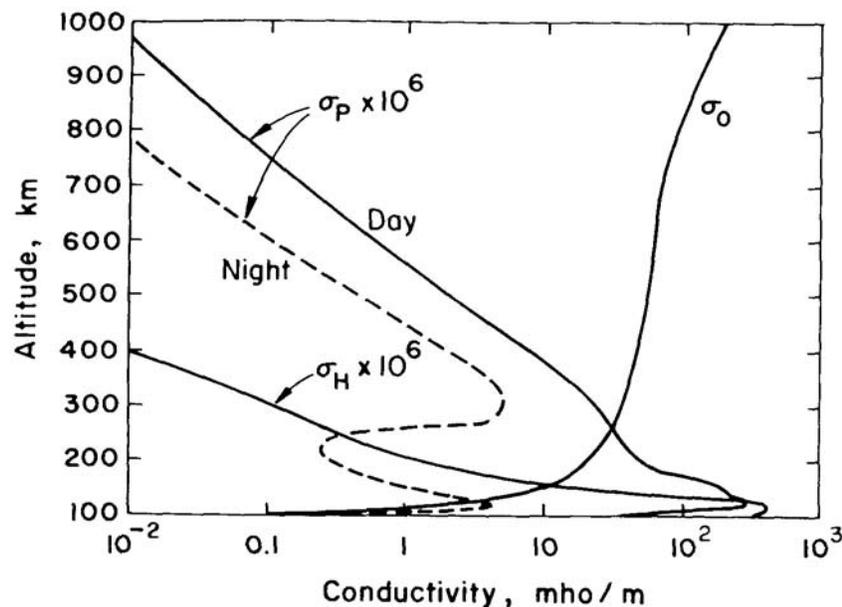
$$J_{\perp} = \sigma_P E_{\perp} - \sigma_H \frac{E_{\perp} \times B}{B}$$

$$J_{\parallel} = \sigma_0 E_{\parallel}$$

$\sigma_P$  = Pedersen conductivity

$\sigma_H$  = Hall conductivity

$\sigma_0$  = parallel conductivity



- Particle precipitation can change  $\sigma$ , so **feedback interaction** with the magnetosphere

# Steady-state Energy Transfer to the Ionosphere: Joule Heating



- Downward Poynting flux from the magnetosphere is dissipated in the ionosphere
- **Joule heating** rate:  $\mathbf{J} \cdot \mathbf{E}$
- Energy dissipated by the current parallel to  $\mathbf{E}_{\perp}$  – **Pedersen current**
- Joule heating dependence on characteristic energy; more energetic particles penetrate more deeply into the ionosphere, where  $\sigma_H > \sigma_P$ , therefore Joule heating is less important for them



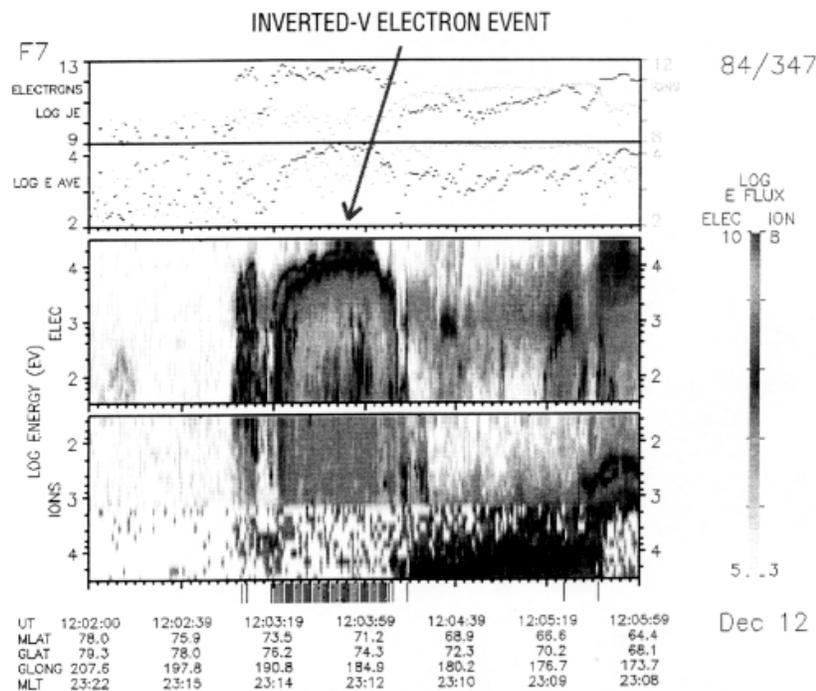
**Vertical structure of conductivity also important**

- Aside: Joule heating can also lead to ionospheric thermal ion outflow events.

# Auroral energization region; electron beams

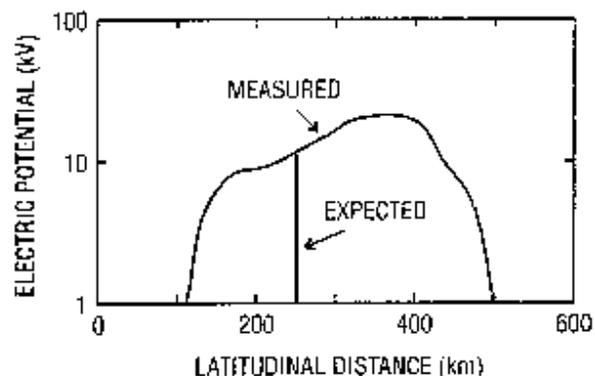


- “Inverted-V” electron events in auroral acceleration region (2000 - 4000 km): energy of electrons indicates the electric potential



- Discrepancy between inverted-V events and discrete auroral arcs

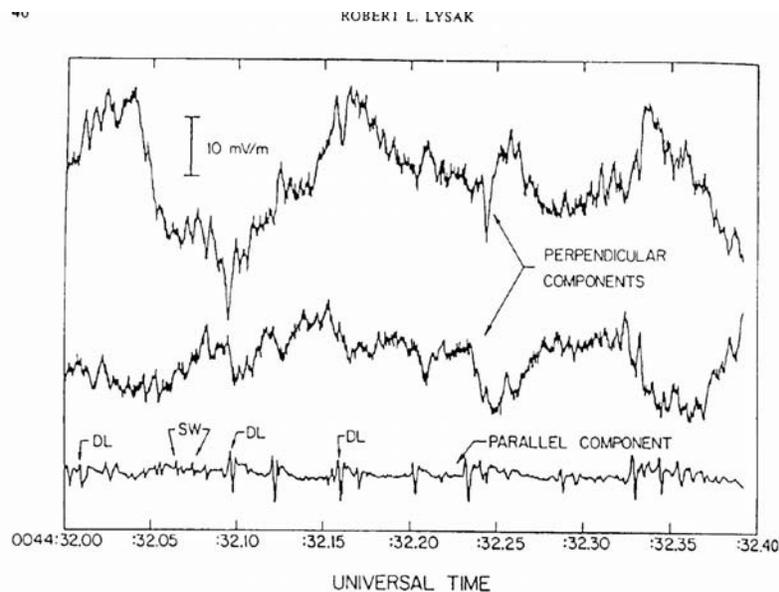
THE LATITUDINAL WIDTH OF AN INVERTED-V ELECTRON EVENT IS OFTEN NEARLY TWO ORDERS OF MAGNITUDE LARGER THAN THE APPARENT THICKNESS OF A THIN DISCRETE ARC.



# $E_{\parallel}$ Formation



- Parallel electron beams (1-10 keV) energized by **parallel potential drop**
- They excite atoms in the neutral atmosphere – **aurora**
- $E_{\parallel}$  theories:
  - Macroscopic: model the macroscopic current
  - Microscopic: how  $E_{\parallel}$  arises from microscopic effects (instabilities)
- Observations:  $E_{\parallel}$  not continuous, but series of localized potential steps: **double layers**

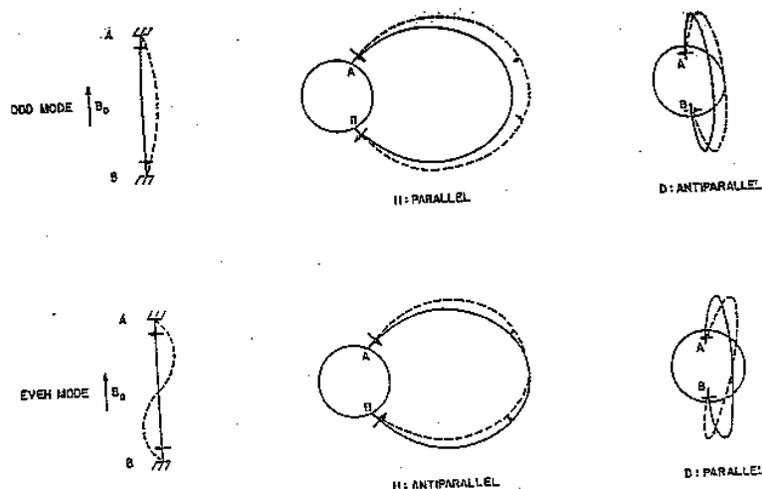


- **Double layer** formation – plasma turbulence effects (nonlinear effects due to trapping of ions and electrons in localized potential bumps)

# Time-dependency: Role of Alfvén Waves



- System **not in steady-state** – inclusion of temporal variation required
- Ionosphere “cold” (low  $\beta$ ); Cold MHD plasma wave theory – 3 wave modes:
  - Fast (magnetosonic) wave
  - Slow wave
  - (shear) Alfvén wave, with  $V_A = B_0/(4\pi\rho)^{1/2}$
- Shear waves exists in the magnetosphere (“natural modes” - **field-line resonances – FLR**) and couple to the ionosphere



# Similarity between static Pedersen current and Alfvén waves



- Physical quantities in the Alfvén wave:

- Fluid velocity: 
$$v = \pm \frac{b}{\sqrt{4\pi\rho}}$$

- Perpendicular electric field: 
$$E_{\perp} = \pm \frac{B_0 \times b}{c \sqrt{4\pi\rho}}$$

- Perpendicular current: 
$$I_{\perp} = \pm \frac{c^2}{4\pi V_A} E_{\perp}$$

(currents defined as  $I_{\perp} = \int dz j_{\perp}$ )

- Analogy:

- Ohm's law with Pedersen  $\Sigma$ : 
$$I_{\perp} = \Sigma_P E_{\perp}$$

- Alfvén “conductivity”: 
$$\Sigma_A = \frac{c^2}{4\pi V_A}$$

- Difficult (but possible) to distinguish between static current patterns (Pedersen) and propagating Alfvén waves

- Conductivities different:  $\Sigma_P > 1 \text{ mho}$

$$\Sigma_A \leq 0.1 \text{ mho}$$

# Alfvén Wave Reflection; Ionospheric Alfvén Resonator (IAR)



- Due to difference between  $\Sigma_A$  and  $\Sigma_P \Rightarrow$  ionospheric **reflection of Alfvén waves**
- By matching currents of incident/reflected wave with ionospheric currents  $\Rightarrow$  reflection coefficient

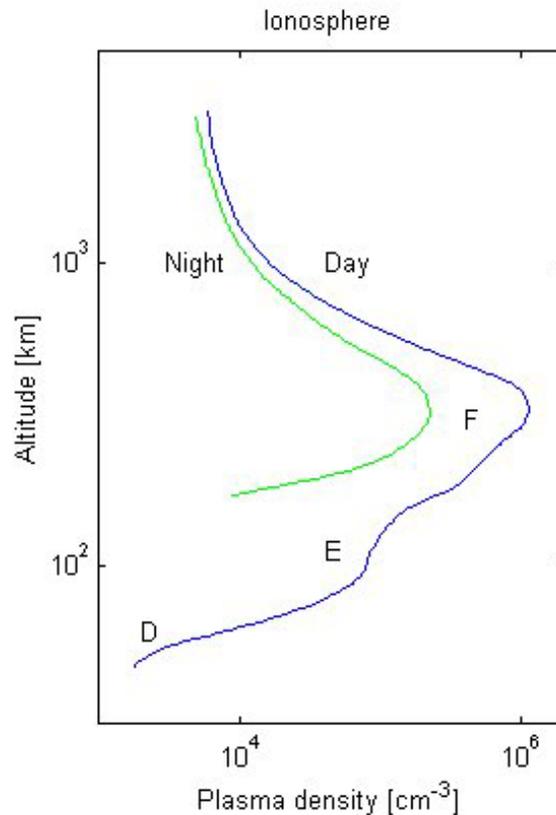
$$R = \frac{\Sigma_A - \Sigma_P}{\Sigma_A + \Sigma_P}$$

- $V_A$  not constant  $\Rightarrow$  waves can be trapped between two altitudes characterized by large  $V_A$  : the lower limit at F-layer, the upper limit at about 3000 km
- Formation of resonant cavity modes – **Alfvén resonator**
  - Frequencies of 0.1 – 1 Hz and higher
  - Can have growing modes - ionospheric feedback instability (beam instability)

# Ionospheric waveguide



High electron density at lower altitudes  
compressional (magnetosonic) waves  
also propagate or resonate –  
ionospheric waveguide in F region of  
the ionosphere for **magnetosonic  
waves**



# Wave-particle interactions



- MHD picture breaks down for large  $k_{\perp}$
- **Inertial term** important at low altitudes ( $v_e < V_A$ )

$$\omega = \frac{k_{\parallel} V_A}{\sqrt{1 + k_{\perp}^2 c^2 / \omega_{pe}^2}}$$

- Kinetic effects (**finite Larmor radius**) important at higher altitudes ( $v_e > V_A$ )

$$\omega = k_{\parallel} V_A \sqrt{1 + k_{\perp}^2 \rho^2}$$

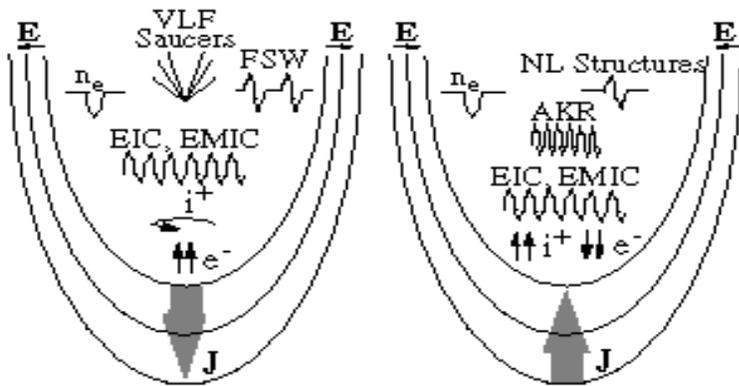
- Both regimes called “**kinetic Alfvén waves**”
- Kinetic Alfvén waves have  $E_{\parallel}$  !
  - Electron population accelerated in bulk – current-driven instabilities may appear
  - Excitation of instabilities by the current in the waves – **formation of steady electric field** through the formation of double layers or anomalous resistivity

# Recent Observations of Auroral Energization Region



- **FAST mission** (Fast Auroral SnapshoT) – launched in 1996

## The Symmetric Auroral Current Regions



- |  |   |
|--|---|
| 1. Downward current region.<br>$\Downarrow J$  | 1. Upward current region.<br>$\Uparrow J$                                     |
| 2. Diverging electrostatic shocks. $\underline{E} \quad \underline{E}$               | 2. Converging electrostatic shocks. $\underline{E} \quad \underline{E}$       |
| 3. Small-scale density cavities.<br>$\frac{n_e}{\downarrow}$                         | 3. Large-scale density cavity.<br>$\frac{n_e}{\downarrow}$                    |
| 4. Up-going, field-aligned electrons. Counter-streaming electrons.<br>$\Uparrow e^-$ | 4. Down-going "inverted-V" and field-aligned electrons.<br>$\Downarrow e^-$   |
| 5. Ion heating transverse to B. Energetic ion conics.<br>$\leftarrow i^+$            | 5. Up-going ion beam. Ion conics.<br>$\Uparrow i^+$                           |
| 6. ELF electric field turbulence. Ion cyclotron waves.<br>                           | 6. Large-amplitude ion cyclotron waves and electric field turbulence.<br>     |
| 7. Fast solitary waves: three-dimensional, rapidly moving electron holes.<br>        | 7. Nonlinear, time-domain structures associated with ion cyclotron waves.<br> |
| 8. VLF saucer source region.<br>   | 8. AKR source region.<br>   |

- Upward and downward currents
- Parallel electric fields
- Particle beams
- Plasma heating
- Host of wave-particle interactions
- Higher-frequency waves (ion-cyclotron)

# Summary



- **Auroral zone – “transmission line”**, carrying EM energy from a magnetospheric generator region to a load region (ionosphere)
  - **Joule dissipation** in the ionosphere
  - Non-linear losses (plasma turbulence)
- Auroral particle energization
  - Particle acceleration region (2000-4000 km)
  - **Parallel potential drop**
- Time-dependent picture
  - **Alfvén waves**
  - Wave-particle interactions:
- M-I system – dynamical system; challenging object of study – rich physics
  - Plasma kinetic theory
  - More global aspects – fluid models
  - Similar current systems and particle energization processes are likely to be present in other astrophysical processes (e.g. solar flares, accretion disks)

# Issues in M-I Coupling Campaign WG2



- **Ionospheric conductance**
  - global distribution ?
  - temporal and spatial variability
- **Auroral Plasma Energization**
  - Relationship between precipitating electron flux and field aligned currents
  - To what extent is MI coupling hemispherically conjugate and synchronous?
  - What processes determine the formation and structure, including length scales, time scales and altitude, of auroral acceleration regions?
- **Multi-scale Processes**
  - Manifestations: discrete aurora, filamentary and layered auroral structures, polar cap arcs
  - how do the different scale sizes interact ?
  - Are averages of energy dissipation meaningful ?  
 $\langle \Sigma \rangle \langle E^2 \rangle \neq \langle \Sigma E^2 \rangle$  How far off are they?
  - What scales contribute most to energy dissipation?)
- **More general question:** Does MI coupling regulate (or how does it regulate) magnetospheric convection, magnetotail dynamics, and solar wind-magnetosphere coupling?
- **Energy Budget Challenge** – the energy flow from the magnetosphere to low altitudes and its myriad pathways for deposition in the ionosphere and lower magnetosphere



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